Measuring CP nature of top-Higgs couplings at the future Large Hadron electron Collider

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We investigate the sensitivity of top-Higgs coupling by considering the associated vertex as CP phase \( \gamma_L \) dependent through the process \( p e^- \rightarrow t H \nu \) in the future Large Hadron electron Collider. In particular the decay modes are taken to be \( h \rightarrow bb \) and \( t \rightarrow \ell \text{ lepton} \) mode. Several distinct \( \gamma_L \) dependent features are demonstrated by considering observables like cross sections, top-quark polarisation, rapidity difference between \( h \) and \( t \) and different angular asymmetries. Luminosity \( (L) \) dependent exclusion limits are obtained for \( \gamma_L \) by considering significance based on fiducial cross sections at different \( \sigma \)-levels. For electron and proton beam-energies of 60 GeV and 7 TeV respectively, at \( L = 100 \text{ fb}^{-1} \), the regions above \( \pi/5 < \gamma_L < \pi \) are excluded at 92\% confidence level, which reflects better sensitivity expected at the Large Hadron Collider. With appropriate error fitting methodology we find that the accuracy of SM top-Higgs coupling could be measured to be \( \kappa = 1.00 \pm 0.17(0.08) \) at \( \sqrt{s} = 1.3(1.8) \text{ TeV} \) for an ultimate \( L = 1 \text{ ab}^{-1} \).

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1. Introduction

The recent discovery of the Higgs boson at the Large Hadron Collider (LHC) serves as the last step in establishing the particle content of the Standard Model (SM). The next step that has been undertaken is the characterisation of its properties regarding spin, CP-nature and the nature of interaction with other particles. While the spin-0 nature of the Higgs boson has been established by the experiments [1–5] and a complete CP-odd nature excluded at a 99.98\% confidence limit (C.L) [6,7], the possibility remains that the Higgs boson could still be an admixture of CP-odd and even states. Investigation of this possibility in a future Large Hadron electron Collider (LHeC) is the goal of this article via a detailed analysis of the associated production of the Higgs boson with an anti-top quark.

Since in the SM the Higgs boson coupling to fermions is directly proportional to the mass of the fermions, the Yukawa coupling associated with the third generation is important in the context of investigating the properties of the Higgs boson. Deviations in the top-Higgs coupling directly affect the production cross section of Higgs boson at the colliders, while changes in the bottom-Higgs coupling affect the total branching ratios.

Here we study the associated production of the Higgs boson with an anti-top quark at the future \( e^- p \) collider which employs a 7 TeV proton beam from a circular \( pp \) collider, and electrons from an Energy Recovery Linac (ERL) being developed for the LHeC [8,9]. The choice of an ERL energy of electron of \( E_e = 60 \) to 120 GeV, with available proton beam energy \( E_p = 7 \text{ TeV} \) provides centre of mass energy of \( \sqrt{s} \approx 1.3 \) to 1.8 TeV. While the LHC is clearly energetically superior, the LHeC configuration is advantageous for the following reasons: (i) since initial states are asymmetric, backward and forward scattering can be disentangled, (ii) it provides a clean environment with suppressed backgrounds from strong interaction processes and free from issues like pile-ups, multiple interactions etc. (iii) such machines are known for high precision measurements of the dynamical properties of the proton allowing simultaneous tests of electroweak and QCD effects. A detailed report on the physics and detector design concepts of the LHeC can be found in the Ref. [8]. A distinguishing feature of the \( e^- p \) collider is that the production of the Higgs is only due to electroweak processes [10,11] and as noted above, since the \( e^- \) and \( p \) energies are different, the machine can also produce interesting patterns of...
kinematic distributions that one can exploit to explore the CP nature of the Higgs boson.

Denoting the CP-odd (CP-even) components of the top-Higgs coupling by $C_T^\pm$ ($C_S^\pm$), the updated bound on the CP top-Higgs couplings by combining the LHC Run-1 and Run-2 Higgs data sets allows the ranges $|C_T^+| < 0.37$ and $0.85 < C_S^+ < 1.20$, which is stronger than the previous LHC Run-1 bound $|C_T^+| < 0.54$ and $0.68 < C_S^+ < 1.20$. We note here that a future precision measurement of the process $e^+e^- \rightarrow h\gamma$ with an accuracy of 0.5% will be able to constrain $|C_T^+| < 0.19$ at a 240 GeV $e^+e^-$ Higgs factory [12]. Various studies on anomalous top-Higgs coupling in associated production of Higgs and top quark can be found in [13–16].

The article is organised as follows: We discuss the formalism by introducing a generalised CP-phase dependent top-Higgs coupling Lagrangian in Section 2. In Section 3 simulation and parton-level analyses of the process emphasising relevant kinematic observables are discussed. Also in this section we provide luminosity dependent exclusion limits of phases corresponding to the top-Higgs coupling. Finally, in Section 4 we conclude with inferences and summary. Though the whole focus of this study is in the LHeC environment, we also discuss and compare our results with those expected at the LHC.

2. Formalism

In the SM, the Yukawa coupling of the third generation of quarks is given by

$$ L_{\text{Yukawa}} = -\frac{m_t}{v} tt h - \frac{m_b}{v} bb h, $$

where $v = (\sqrt{2} G_F)^{-1/2} = 2 m_W / g \simeq 246$ GeV, and $m_t$ ($m_b$) is the mass of the top (bottom) quark. Due to the pure scalar nature of the Higgs boson in the SM, here the top- and bottom-Higgs couplings are completely CP-even. To investigate any beyond the SM (BSM) nature of the Higgs-boson as a mixture of CP-even and CP-odd states, we write a CP-phase dependent generalised Lagrangian as follows [17]:

$$ L = \frac{-m_t}{v} [k \cos \zeta_t + i \gamma_5 \sin \zeta_t] t \bar{h} - \frac{m_b}{v} [\cos \zeta_b + i \gamma_5 \sin \zeta_b] b \bar{h}. $$

(2)

Here $\zeta_t$ and $\zeta_b$ are the phases of the top-Higgs and bottom-Higgs couplings respectively. It is clear from the Lagrangian in Eq. (2) that $\zeta_{t,b} = 0$ or $\zeta_{t,b} = \pi$ correspond to a pure scalar state while $\zeta_{t,b} = \frac{\pi}{2}$ to a pure pseudo scalar state. Thus, the ranges $0 < \zeta_{t,b} < \pi/2$ or $\pi/2 < \zeta_{t,b} < \pi$ represent a mixture of the different CP-states. The case $\kappa = 1$, $\zeta_t = 0$ corresponds to the SM. In terms of $C_T^S$ and $C_S^T$, we can also translate $\zeta_t = \tan^{-1}(C_T^S / C_T^T)$.

At the LHeC, the top-Higgs couplings can be probed via associated production of Higgs-boson with anti-top quark $p e^- \rightarrow t \bar{h} \nu_e$ - it is thus necessary to consider a 5-flavour proton including the b-quark parton distribution. The Feynman diagrams for the process are shown in Fig. 1. It is important to notice that in this process three important couplings are involved, namely $hWW$, $Wtb$ and the top-Higgs ($ttH$). A detailed study of $hWW$ and $Wtb$ couplings at the $e^+e^-$ collider has been performed in Refs. [11,18] and [19], respectively. For our studies we do not consider the BSM bottom-Higgs coupling since the effect of the phase $\zeta_b$ on the total production cross section or kinematics of top-Higgs production at the LHeC is negligible. Thus in what follows, we simply set $\zeta_b = 0$.

As noted in Ref. [17] in the context of the LHC, quantitatively an interesting feature can be observed: in the pure SM case there is constructive interference between the diagrams shown in Fig. 1a and Fig. 1c for $\zeta_t > \pi/2$ resulting in an enhancement in the total production cross section of associated top-Higgs significantly. This is also true for $\zeta_t < \pi/2$ - however the degree of enhancement is much smaller owing to the flipped sign of the CP-even part of the coupling.

3. Simulation and analysis

We begin our study to probe the sensitivity of the top-Higgs couplings in terms of $\zeta_t$ by building a model file for the Lagrangian in Eq. (2) using FeynRules [20], and then simulating the charged current associated top-Higgs production channel $p e^- \rightarrow t \bar{h} \nu_e$ (see Fig. 1), with $h$ further decaying into a $b \bar{b}$ pair and the $t$ decaying leptonically in the LHeC set-up with centre of mass energy of $\sqrt{s} \approx 1.3$ TeV. In this article we perform the analysis at parton level only where for signal and background event generation we use the Monte Carlo event generator package MadGraph5 [21]. We use NNPDF23LO1 [22,23] parton distribution functions for all event generations. The factorisation and renormalisation scales for the signal simulation are fixed at $\mu_F = \mu_R = (m_t + m_b)/4$ while background simulations are done with the default MadGraph5 [21] dynamic scales. The $e^-$ polarisation is assumed to be $-80\%$. We now list and explain various kinematic observables that can serve as possible discriminants of a CP-odd $t \bar{h}$ coupling.

3.1. Cross section studies

In Fig. 2, we present the variation of the total cross section against the electron beam energy for the signal process.
Table 1

Cross sections of signal and backgrounds in charged current (cc), neutral current (nc) and photo-production (pho) modes for \( E_L = 60 \text{ GeV} \) and \( E_p = 7 \text{ TeV} \) as explained in the text. Here \( X \) could be either of missing energy or electron and \( j \) is all possible combinations of light-, c- and b-quarks and gluons. For this estimation we use basic cuts as mentioned in text and electron polarisation is taken to be \(-0.8\).

<table>
<thead>
<tr>
<th>Process</th>
<th>cc (fb)</th>
<th>nc (fb)</th>
<th>pho (fb)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Signal:</td>
<td>( 1.98 \times 10^{-2} )</td>
<td>-</td>
<td>-</td>
</tr>
<tr>
<td>( Wjjj + X, \not!{h} )</td>
<td>( 2.05 \times 10^{-2} )</td>
<td>( 3.18 \times 10^{-1} )</td>
<td>( 3.40 \times 10^{-1} )</td>
</tr>
<tr>
<td>( Wjjj + X, \not!{t} )</td>
<td>( 4.18 \times 10^{-1} )</td>
<td>( 3.16 \times 10^{-1} )</td>
<td>( 3.41 \times 10^{-1} )</td>
</tr>
<tr>
<td>( Wjjj + X, \not!{h} )</td>
<td>( 4.16 \times 10^{-1} )</td>
<td>( 3.18 \times 10^{-1} )</td>
<td>( 3.41 \times 10^{-1} )</td>
</tr>
</tbody>
</table>

\[ p e^- \rightarrow t \bar{h} n_{\nu_e}, \] by considering un-polarised and \(-80\%\) polarised \( e^- \) beam. Also, the effect of branchings of \( h \rightarrow b \bar{b} \) and the \( t \) decay for both leptonic and hadronic modes are shown. Possible background events typically arise from \( W^+ \) multi-jet events, \( Wbb \) with missing energy which comes by considering only top-line \( \not\!{t} \), only Higgs-line \( \not\!{h} \) and without top- and Higgs-line \( \not\!{t} \) in charged and neutral current deep-inelastic scattering and in photo-production by further decaying \( W \) into leptonic mode. In Table 1 we give an estimation of cross sections for signal and all possible backgrounds imposing only basic cuts on rapidity \( |y| \leq 10 \) for light-jets, leptons and \( b \)-tagged jets, the transverse momentum cut \( p_T \geq 10 \text{ GeV} \) and \( \Delta \eta_{\text{min}}^1 \) = 0.4 for all particles.

We now estimate the sensitivity of the associated top-Higgs production cross-section, \( \sigma(\xi_t) \), as a function of the CP phase of the \( t \)th-coupling as shown in Fig. 3 by considering \( E_L = 60 \) and \( 120 \) GeV with fixed \( E_p = 7 \) TeV. The scale uncertainties are taken as \( (m_t + m_h)/8 \leq \mu_s = \mu_R = (m_t + m_h)/2 \). Here \( \sigma(\xi_t = 0) \) corresponds to the SM cross section. We notice that the cross section is very sensitive to \( \xi_t \) in the region \( \xi_t > \frac{\pi}{2} \) where the interference between the diagrams becomes constructive. Below \( \xi_t = \frac{\pi}{2} \) the interference is still constructive though its degree decreases with \( \xi_t \), thus increasing the cross section by around 500\% at \( \xi_t = \frac{\pi}{2} \) which corresponds to the pure CP-odd case. On the other hand, for pure CP-even case \( \xi_t = \pi \) with opposite-sign of \( t \)th-coupling the cross section can be enhanced by up to 2400\% for \( E_L = 60 \) GeV. Notice that for the case \( E_L = 120 \) GeV, \( \sigma(\xi_t) \) displays a similar shape with enhanced cross sections with respect to \( E_L = 60 \) GeV case. The scale uncertainty on an average is approximately 7(9)\% for \( E_L = 60(120) \) GeV in the whole range of \( \xi_t \).

\[ \Delta \eta_{\text{min}}^1 = \sqrt{(\Delta \phi)^2 + (\Delta \eta)^2}, \] where \( \phi \) and \( \eta \) are the azimuthal angle and rapidity respectively of particles into consideration.

3.2. Rapidity difference between the anti-top and the Higgs

In Refs. [11,18] it was suggested that in order to explore the tensorial spin-CP nature of \( hhW^+W^- \) and \( hhW^+W^- \) vertices, azimuthal angle correlation between missing energy and scattered jets is a good observable. Also further studying the asymmetry based on such observables proves to be an excellent tool for any BSM nature of the associated couplings. Here and in the next subsections we include such observables in our studies with different combinations of final state particles as a function of \( \xi_t \). We begin with the sensitivity of BSM aspects of the \( t \)th coupling in the rapidity difference between the anti-top quark and the Higgs boson distribution, \( \Delta y_{\text{ht}} \).

In Fig. 4 we present the normalised \( \Delta y_{\text{ht}} \) distribution for a few chosen values of \( \xi_t \). Any BSM physics effect can be observed by comparing the shape corresponding to the SM case \( \xi_t = 0 \). We find that the distribution features for the different values of CP phase split into two distinguishable regions when \( \Delta y_{\text{ht}} < 1 \) and \( 1 < \Delta y_{\text{ht}} < 3 \). In the former, most values of \( \xi_t \) are seen to correspond to distributions larger than the SM case, while the second region presents a complementary behaviour. The distortion in the shape for \( \xi_t > 0 \) is the effect of mixing between CP-even and odd components of the \( t \)th vertex following the Lagrangian in Eq. (2). Overall, with the inclusion of spin-1\(^+\) BSM admixture, the \( \Delta y_{\text{ht}} \) distribution is pushed towards lower values and acts as a potential discriminator to explore the CP-nature of \( t \)th-coupling. Similar studies are used to probe the tensor structure of \( hVV \) (\( V = W^\pm, Z \)) coupling at the LHC and one such study of the Higgs boson production in the vector boson fusion mode is performed in [27] by taking the rapidity difference between the Higgs and the leading parton.

\[ \mu_s = \frac{r_{\text{observed}}}{r_{\text{SM}}} \]

Note that at the LHC the production of associated Higgs boson with top-quark is possible via double and single-top quarks and is different from LHeC where the environment and centre of mass energies are different. The signal strength is defined as \( \mu_s = \frac{r_{\text{observed}}}{r_{\text{SM}}} \).
3.3. Top quark polarisation

The large top-quark mass $m_t = 172.84 \pm 0.70$ GeV [28] indicates that the top could potentially play a singular role in the understanding of electroweak symmetry breaking in BSM scenarios. Since the decay width of the top exceeds $\Lambda_{QCD}$, the top decays before hadronising and thus its spin information is preserved in the differential distribution of its decay products. With the Higgs coupling to top modified, it is reasonable to expect an asymmetry in the production of tops of different polarisations and the effect of $\zeta_t$ on this asymmetry.

We define the degree of longitudinal polarisation $P_t$ of the top quark as

$$P_t = \frac{N_+ - N_-}{N_+ + N_-} = \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-},$$

where $N_+$ and $N_-$ denote the number of events with positive and negative helicity anti-top quarks respectively, which can be rewritten in terms of the corresponding cross sections $\sigma_{\pm}$ in Fig. 5. We present $P_t$ in the process $p e^- \to t \bar{t} \nu_e$ at the LHeC as a function of $\zeta_t$. We obtain $N_\pm$ or $\sigma_\pm$ using the helicity amplitudes in MadGraph5. It can be seen from the plot that the degree of polarisation is quite sensitive over the entire range of $\zeta_t$ since the CP-odd coupling violates parity for any non-zero $\zeta_t$. It is interesting to note that if Fig. 1c is the only diagram that contributed to $P_t$, then the fraction of right-handedly polarized anti-top quark would increase as $\zeta_t$ increases from 0 and reach a maximum at $\zeta_t = \pi/2$ and then fall. However, the presence of other diagrams means that the plot is not symmetric about $\zeta_t = \pi/2$. The general features of $P_t$ in Fig. 5 can be understood as the effect of interference among the diagrams in Fig. 1a, Fig. 1b (from where right-handed anti-top quarks are produced) and the Higgs-bremsstrahlung diagram Fig. 1c, which contains the CP-violating sin $\zeta_t$.

As mentioned before, information of the spin of the top is preserved in its decay products and the angular distribution of its decay products can be parametrised as:

$$\frac{1}{\Gamma_f} \frac{d\Gamma_f}{d\cos \theta_f} = \frac{1}{2} (1 + \alpha_f P_t \cos \theta_f),$$

where $f$ is the type of top decay product, $\theta_f$ is the angle between $f$ and the top-quark spin quantisation axis measured in the rest frame of the top-quark, and $\Gamma_f$ denotes the partial decay width corresponding to $f$. For the decay mode $t \to b + W^\pm \to b + l^\pm + \nu_l$, at lowest order, $\alpha_{W} = -\alpha_b = 0.39$, $\alpha_\ell = -0.3$, $\alpha_1 = 0$ [29], with small QCD corrections to these values [30,31]. The charged lepton $l^\pm$ (or the down-type quark $d$ in a hadronic decay of the intermediate $W$) is nearly 100% correlated with the top quark spin which means that the $l^\pm$ or $d$ is much more likely to be emitted in the direction of the top quark spin than in the opposite direction. It is a well known fact that the energy and momentum of leptons can be measured with high precision at the LHC and the same is true for the LHeC as well, so we focus on the leptonic decay mode of the anti-top for asymmetries in angular observables studies in what follows.

3.4. Cut-based event optimisation

Before discussing the angular observables for this study, it is important to discuss the optimisation of SM signal and background events as mentioned in Section 3.1. Angular observables are affected due to kinematic cuts and hence it is better to analyse events after optimising the signal with respect to backgrounds. The full SM signal process for this analysis is $p e^- \to t \bar{t} h \nu_e$, with $h \to b \bar{b}$ and $t \to W^- b$, $W^- \to l^- \nu_l$ ($l^\pm = e^\pm, \mu^\pm$). After preliminary analysis of various kinematic distributions of final state particles of the SM signal and all possible leptonic backgrounds, we employ the following criteria to select events: (i) $p_T \geq 20$ GeV for $b$-tagged jets and light-jets, and $p_T \geq 10$ GeV for leptons. (ii) Since the LHeC collider is asymmetric, event statistics of final state particles are mostly accumulated on the left or right sides of the transverse plane $\eta = 0$ (depending on the initial direction of $p$ and $e^-$) - we select events within $-2 \leq \eta \leq 5$ for $b$-tagged jets while $2 \leq \eta \leq 5$ for leptons and light-jets. (iii) The separation distance of all final state particles are taken to be $\Delta R > 0.4$. (iv) Missing transverse energy $E_T > 10$ GeV to select the top events. (v) Invariant mass windows for the Higgs through $b$-tagged jets and the top are required to be $115 < m_{bb} < 130$ GeV and $160 < m_t < 177$ GeV respectively which are important to reduce the background events substantially. In these selections the $b$-tagging efficiency is assumed to be 70%, with fake rates from c-initiated jets and light jets to the $b$-jets to be 10% and 1% respectively. These constitute our event selection criteria which we use in the subsequent analysis.

There are two major difficulties in reconstructing the Higgs boson and the top in the process $p e^- \to t \bar{t} h \nu_e \to (W^- b) h \nu_e \to l^- \nu_l b b \nu_l \nu_e$: (a) Choosing appropriate $b$-tagged jets – in the final state we have 3 $b$-tagged jets with two originating from $h$ decay and one from the decay of $t$ and (b) the source of missing energy comes from both the production process and from $W^\pm$ decay. Since we performed parton-level analysis, we read the event files generated from the Monte Carlo generator and by reading appropriate identities we obtained information about the origin of $b$-tagged jets and neutrino and the corresponding four-momenta information was used for the analysis. Although the detector-level analysis is beyond the scope of this article, we mention briefly that for distinguishability of $b$-jets the solution is to take into account the $p_T$ ordering of all $b$-tagged jets and since top-quark is heavier than the Higgs boson, the leading-$p_T$ $b$-jet can be identified as the decay product of top-quark, and the sub-leading and next to sub-leading $p_T$-ordered $b$-jets can be used to reconstruct Higgs boson.

To reconstruct the top, substantial requirement on missing energy and top-quark invariant mass formula $m_T^2 = (m_T + m_{b1})$ can be used, where $m_T$ is transverse mass observable to reconstruct $W$-boson and $m_{b1}$ is the mass of leading-$b$ jet and is given as:

$$m_T = \sqrt{2 p_T^1 p_T^2 (1 - \cos(\phi_1 - \phi_2))},$$

where $\cos(\phi_1 - \phi_2)$ is the angle between the electron and neutrino in the transverse plane, and $\phi_1$ ($\phi_2$) is the azimuthal angle of the electron (neutrino). However, it is to be noted that $m_T$ is also inefficient when there are more than one source of missing energy and hence alternative method should be explored.
3.5. Angular observables in terms of asymmetries

After this short discussion on event selection criteria, we now discuss observables based on angular asymmetry between different final state particles. We construct the asymmetry from the differential distribution of kinematic observables using the final leptons and b-taggerged jets. These asymmetries are studied only for signal processes as a function of \( \xi_t \). The angular asymmetries with respect to polar angle \( \cos \theta^i \) and the azimuthal angle difference \( \Delta \phi_{ij} \) are defined to be:

\[
A_{\theta^i j} = \frac{N^2(\cos \theta^i_j > 0) - N^2(\cos \theta^i_j < 0)}{N^2(\cos \theta^i_j > 0) + N^2(\cos \theta^i_j < 0)},
\]

\[
A_{\Delta \phi_{ij}} = \frac{N^A(\Delta \phi_{ij} > \pi/2) - N^A(\Delta \phi_{ij} < \pi/2)}{N^A(\Delta \phi_{ij} > \pi/2) + N^A(\Delta \phi_{ij} < \pi/2)},
\]

where \( i \) and \( j \) are any two different final state particles. Using binomial distribution we use the following formula to calculate the statistical uncertainty (\( \delta_{A_{\alpha}} \)) in the measurement of these asymmetries (\( A_{\alpha} \)):

\[
\delta_{A_{\alpha}} = \sqrt{\frac{1 - A^2_{\alpha} (\xi_t)}{\sigma_{\alpha}^2 \cdot L}}, \quad (\alpha = \theta^i, \Delta \phi_{ij})
\]

where \( \sigma_{\alpha} \) is the total cross section of signal events as a function of \( \xi_t \) and \( L \) is the total integrated luminosity.

In Fig. 6, we show the asymmetries between the charged lepton and the b from \( f \) decay (denoted by \( b_1 \) in the plot) as functions of \( \xi_t \). We can see that the asymmetries in \( \Delta \phi_{t,ij} \) and \( \cos \theta^i \) follow the top polarisation curve to some extent in that they fall till \( \xi_t \approx \pi/4 \). We find that beyond \( \xi_t \approx \pi/2 \), the curves flatten. As explained in the Section 3.3 the shapes in these asymmetry observables are also influenced by interference among the Feynman diagrams shown in Fig. 1. Overall we can conclude that these asymmetry observables can serve as good discriminators for a non-zero \( \xi_t \), particularly for \( \xi_t < \pi/2 \) where the difference from the \( \xi_t = 0 \) case is more pronounced.

3.6. Exclusion limits

In Section 3.5 we observed that asymmetry observables based on differential distributions of \( \cos \theta^i \) and \( \Delta \phi_{t,ij} \) show distinct features in terms of shape although quantitatively not very sensitive. Therefore we construct another asymmetry observable by considering the polar angle between the sub-leading b-taggerged jet and the lepton from W⁻ decay, i.e., \( \cos \theta^j \), which is comparatively more sensitive (quantitatively). In Fig. 7, we show the asymmetry \( A_{\cos \theta^j} \) as a function of \( \xi_t \) for \( E_x = 60 \) and 120 GeV with \( E_p = 7 \) TeV. The statistical uncertainties are calculated using the formula in Eq. (7) for \( \xi_t = 0 \) and explicitly given as:

\[
\delta A_{\cos \theta^j} = \sqrt{\frac{1 - (A_{\cos \theta^j}^{SM})^2}{\sigma_{SM} \cdot L}},
\]

where \( \sigma_{SM} \) is total cross section of the SM signal and \( A_{\cos \theta^j}^{SM} \) is numerical value of corresponding SM asymmetry. Therefore at the luminosity of \( L = 1 \) ab⁻¹, \( A_{\cos \theta^j} \) used to determine \( \xi_t \) within \( \pi/3 \) and \( 3\pi/5 \) (\( \pi/6 \) and \( 3\pi/10 \)) at 1σ and 2σ C.L. respectively for \( E_x = 60 \) (120) GeV. This indicates that at low \( L \) the sensitivity tends to be poorer than this, so next we use fiducial inclusive cross sections as another observable to find the exclusion limits.

Based on selection criteria of signal and background events discussed in Section 3.4, we estimated the exclusion regions of \( \xi_t \) as a function of \( L \) in fb⁻¹. The exclusion is based on significance using the Poisson formula \( S = \sqrt{2[(S + B)\log(1 + S/B) - S]} \), where \( S \) and \( B \) are the number of expected signal and background events at a particular luminosity respectively. Here we used 10σ systematic uncertainty for background yields only. In Fig. 8, we present exclusion contours at various confidence levels for \( E_x = 60 \) GeV – understandably, higher \( \sigma \)-contours demand larger luminosities. It is also seen that there is a kink around \( \xi_t = \pi/2 \) such that for the region \( 0 < \xi_t < \pi/2 \), we need larger luminosities for exclusion. This is in keeping with the feature exhibited in Fig. 3 where the constructive interference between the signal diagrams enhances the cross-section over the SM value much more for \( \xi_t > \pi/2 \) than requiring less luminosity to probe that region. For \( L = 100 \) fb⁻¹, regions above \( \pi/5 < \xi_t \leq \pi \) and \( 3\pi/10 < \xi_t \leq \pi \) are excluded at 2σ and 3σ C.L. While around \( L = 400 \) fb⁻¹ , regions above \( \pi/6 < \xi_t \leq \pi \) and \( \pi/4 < \xi_t \leq \pi \) are excluded at 4σ and 5σ C.L. respectively.

For higher \( E_x = 120 \) GeV, the cross section for signal (background) is enhanced approximately by a factor of 4 (3) and hence the luminosity required for exclusion is smaller compared to the
rapideities of anti-top quark and Higgs-boson $\Delta y_{t\bar{t}}$, and anti-top polarisation $P_1$ show unique features that are distinct from the pure scalar type couplings.

Considering the leptonic decay mode of the anti-top quark and $h \to bb$, we constructed the asymptotic observables $\Delta \phi_{t\bar{t}}$ and $\cos \theta_{t\bar{t}}$. We find that while these show deviations from the SM case in the region $0 < \zeta_t \leq \pi/2$, the curves flatten out beyond that point. This prompted us to construct yet another observable $\cos \phi_{t\bar{t}}$ whose variation with $\zeta_t$ is significant in the entire range $0 < \zeta_t \leq \pi$.

Somewhat counterintuitively, exclusion regions for $\zeta_t$ obtained through fiducial cross section considerations result in better limits than those using asymptotic measurements. Quite strikingly, we find that almost all values of $\zeta_t$ can be excluded at $2\sigma$ ($4\sigma$ ) with an integrated luminosity of $200 \text{ fb}^{-1}$ ($700 \text{ fb}^{-1}$) – these limits are superior to those found in studies at the HL-LHC. While the limits would possibly worsen when one does a full detector level simulation, our analysis gives excellent early signs for the efficacy of the LHeC for coupling measurements.

We conclude that a study of cross-section measurements combined with accurate measurements of kinematic observables can be a powerful probe at the LHeC to uncover the finer details of the nature of the top-Higgs coupling and hope that this study adds to the physics goals of future $e^+ e^-$ colliders.

As mentioned in Section 2, apart from $tth$ coupling the process considered in this study involves $hWW$ and $Wtb$ couplings as well where non-standard anomalous contributions are not negligible – these are studied in Refs. [11,18] and [19] respectively. Since the gauge-scalar ($WWW$) and gauge-fermion ($Wtb$) anomalous couplings involve momentum dependent couplings, the differential distribution of final state particles is affected differently via such effects and can thus be used as an effective discriminant to disentangle the effects of different new physics contributions to the process under investigation. For future studies, a global analysis involving all anomalous non-standard couplings together will be helpful to investigate the potential of precision measurement capabilities of collider facilities like the LHeC.

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